Structures of $d^4$ MH$_3$X: a Computational Study of the Influence of the Metal and the Ligands

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Abstract

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ABSTRACT: Density functional theory (DFT, PBE0, and range separated DFT, RSH + MP2) and coupled-cluster with single and double and perturbative triple excitations (CCSD(T)) calculations have been used to probe the structural preference of d^4 MH_3X (M = Ru, Os, Rh^+, Ir^+, and Re^+; X = H, F, CH_3, CF_3, SiH_3, and SiF_3) and of MX_4 (M = Ru; X = H, F, CH_3, CF_3, SiH_3, and SiF_3). Landis et al. have shown that complexes in which the metal is sd^3 hybridized have tetrahedral and non-tetrahedral structures with shapes of an umbrella or a 4-legged piano stool. In this article, the influence of the metal and ligands on the energies of the three isomeric structures of d^4 MH_3X and MX_4 is established and rationalized. Fluoride and alkyl ligands stabilize the tetrahedral relative to non-tetrahedral structures while hydride and silyl ligands stabilize the non-tetrahedral structures. For given ligands and charge, 4d metal favors more the non-tetrahedral structures than 5d metals. A positive charge increases the preference for the non-tetrahedral structures while a negative charge increases the preference for the tetrahedral structure. The factors that determine these energy patterns are discussed by means of a molecular orbital analysis, based on Extended Hückel (EHT) calculations, and by means of Natural Bond Orbital (NBO) analyses of charges and resonance structures (NRT analysis). These analyses show the presence of through-space interactions in the non-tetrahedral structures that can be sufficiently stabilizing, for specific metals and ligands, to stabilize the non-tetrahedral structures relative to the tetrahedral isomer.

INTRODUCTION

Unsaturated transition metal complexes are a model of reactive intermediates, and their structures and electronic properties can provide useful information on how they could react. These species are often very reactive and hard to characterize by experimental means. Therefore, structure and bonding analyses and computational studies have been useful to establish the structures of these complexes. The structure of unsaturated species does not always follow the rules that apply for saturated systems. For instance, it has been found that the structure of unsaturated species cannot be always predicted by the valence shell electron pair repulsion (VSEPR) theory^1 or by a simple ionic model. WMe_6 and WH_6 are notable cases for the failure of the VSEPR rule. WCl_6 is octahedral^4 while WMe_6 and WH_6 are not. Analysis of the metal–ligand bonding within such complexes shows that the electron delocalization between the π-donor ligand and the formally empty metal d orbital increases the electron density at the metal which, in turn, becomes less unsaturated. Therefore, the structures of polyhydride or polyalkyl complexes are considered as informative of the fundamental structural preferences associated with the metal coordination number and the number of d electrons. Unfortunately, few systems of this type have been characterized experimentally, a clear consequence of the poor stability of highly unsaturated species. For this reason, computational studies have been widely used.5

Landis et al. have shown that valence bond concepts can be used for describing the shape of a variety of hydride and alkyl systems. A primary feature of their studies is that the valence empty (n + 1)p orbitals of a nd metal do not contribute to the metal–ligand bonds, which are constructed from the (n + 1)s and available nd orbitals. Consequently, the shape of the hydride complex, where the metal has λ nonbonding electron pairs, is determined by the optimal spatial distribution of the sd^n (ω = 5 – λ) hybrids set, that is, the spatial distribution that minimizes the interactions between the hybrid orbitals. In general, there is more than one optimal spatial distribution for a given sd^n set and quantum calculations are used to determine the relative energies of these isomeric structures.7a,8,9 Similar results were obtained using the orbitally ranked symmetry...
analysis method (ORSAM) since for low coordinated complexes the metal orbital uses also the \((n + 1)s\) and nd orbitals and the ranking of the structures is also done by quantum calculations.\(^1\) In these studies, only hydride complexes were calculated but other ligands were not explicitly considered. Likewise the influence of the metal was not explored.

Well-defined experimental species are rather limited since the only known complexes of this type are the homoleptic neutral \((M = \text{Ru and Os})\) and cationic \((M = \text{Ir}^+)\) \(\text{d}_4\) \(\text{MR}_4\) \((\text{R} = \text{alkyl and aryl})\) complexes synthesized by Wilkinson. These complexes, \(\text{Ru(C}_6\text{H}_5)_4\),\(^1\) \(\text{Ru(o-Tolyl)}_4\),\(^1\) \(\text{Ru(Mesityl)}_4\),\(^1\) \(\text{Os(C}_6\text{H}_5)_4\),\(^1\) \(\text{Os(o-Tolyl)}_4\),\(^1\) and \(\text{Ir(Mesityl)}_4\),\(^1\) are stable and have a tetrahedral coordination at the metal. However, the bonding analysis of Landis and the ORSAM analysis of Hall predict that other isomers with non-tetrahedral structures should be possible, and calculations for \(\text{RuH}_4\) and \(\text{OsH}_4\) have shown that non-tetrahedral structures could even be preferred or at least competitive, but no rationale for these results is available. We were intrigued by these results, and we wanted to explore more widely the structural preferences of \(\text{d}^4\) \(\text{MH}_4\) and \(\text{MX}_4\) for various metals, charges, and ligands. We present a study based on density functional theory (DFT) calculations, validated by ab initio calculations, for neutral \(\text{MH}_4\) \((\text{M} = \text{Ru and Os}; \text{X} = \text{H, F, CF}_3, \text{CH}_3, \text{SiF}_3, \text{and SiH}_3)\), cationic \(\text{MH}_4^+\) \((\text{M} = \text{Rh and Ir})\), and anionic \(\text{REH}_3\text{X}^-\) \((\text{M} = \text{Rh and Ir})\) complexes. We also present results for homoleptic \(\text{RuH}_4\) \((\text{M} = \text{H, F, CF}_3, \text{CH}_3, \text{SiH}_3, \text{SiF}_3)\) complexes.

### COMPUTATIONAL METHODS

Two methods (1 and 2) were used to obtain the geometries and energies at the DFT level. Additional methods were used to increase the accuracy of the relative energies of the isomers, methods 3 and 4.

**Method 1.** The geometry optimizations were performed at the DFT level of theory with Gaussian 09\(^1\) using the B3PW91 functional.\(^2\) The metal centers were described with the Stuttgart/Dresden (SDD) quasi-relativistic effective core potentials (RECP) and the associated basis sets.\(^3\) The 6-31G(d,p) basis set were used for the other atoms.\(^4\)

**Method 2.** The geometries obtained with the method 1 were used as initial input for optimization with the PBE0 functional\(^5\) and triple-\(\zeta\) quality basis sets on all atoms, de2-TZVPP,\(^6\) with the corresponding RECP for the metal.\(^7\) With rare exceptions, the optimized structures obtained with method 2, and method 1 are similar. Bond lengths only differ by about 0.01 Å, and bond angles differ by less than 5°. The structures and energies obtained with method 2 are used for describing the results in this Article. All stationary points were fully characterized by analytical frequency calculations as either a minimum or a transition state. Intrinsic reaction coordinate (IRC) calculations followed by geometry optimizations were used to associate a transition state to the corresponding two minima. The energies reported in this article do not include the ZPE correction, but it has been verified that the ZPE correction does not modify significantly the relative energies of the isomers. The optimized geometries for all stationary points and the ZPE corrected energies are given in the Supporting Information.

**Method 3.** Single-point short-range (sr) DFT/long-range (lr) ab initio calculations on the DFT (PBE0/de2-TZVPP) geometries were carried out for all structures using the Molpro 2010.1 program.\(^8\) The short-range exchange-correlation terms were treated by the \(\text{sr-PBE}\) variant functional.\(^9\) The long-range portion of the exchange energy is explicitly treated (HF). This step defines a “range-separated hybrid” (RSH) scheme, which is corrected in a second step for the long-range correlation effects by a second order perturbation theory, leading to MP2-like correction. This method is referred to as RSH+MP2.\(^10\) The range-separation parameter \(\mu\) has been set to 0.4 \(\text{a}_0^{-1}\). For these single-point calculations, all atoms were described with quadruple-\(\zeta\) quality basis sets, de2-QZVPP\(^11\) with the corresponding RECP for the metal.\(^12\)

**Method 4.** Single-point coupled-cluster with single and double and perturbative triple excitations (CCSD(T)) calculations on the DFT (PBE0/de2-TZVPP) optimized geometries were carried out for \(\text{MH}_4\) \((\text{M} = \text{Ru and Os}; \text{X} = \text{H, F, CF}_3, \text{CH}_3, \text{SiF}_3, \text{and SiH}_3)\) using the Orca 2.9 program.\(^13\) In this case, all atoms were described with a de2-QZVPP basis sets with the corresponding RECP for the metal.\(^14\)

**Method 5.** Single-point MR-CI and MR-ACPF-2a calculations on the DFT (PBE0/de2-TZVPP) optimized geometries were carried out for \(\text{RuH}_4\) using the Orca 2.9 program.\(^15\) In this case, all atoms were described with a de2-QZVPP basis sets with the corresponding RECP for the metal.\(^16\) Resolution of the identity (RI) approximation was used for the integral transformation with de2-QZVPP/C auxiliary basis sets.\(^17\) For these MRCl and MR-ACPF-2a calculations, the reference space was built from full valence CASSCF orbitals (12 electrons, 10 orbitals) using a selection threshold (Tpre) set to 10⁻⁶. The selection of the excited configurations from the reference space was done with a selection threshold (Tsel) of 10⁻⁷. The effect of the rejected configurations was estimated using second order perturbation theory.

### RESULTS

In \(\text{d}^4\) \(\text{ML}_4\) species, where two \(\text{d}\) orbitals are used to host the four metal electrons, the \((n + 1)s\) and three \(\text{nd}\) metal orbitals are available to form the four metal–ligand bonds. This leads to \(\text{sd}\) hybridization at the metal, and the shape of the complex is determined by the optimal spatial distribution of these four \(\text{sd}\) hybrids. The neutral polyhydride complex, \(\text{d}^4\) \(\text{MH}_4\), is an ideal system in that the \(\text{H}–\text{M}–\text{H}\) angles are very close to those that minimize the overlap between the metal hybrid orbitals at the metal which are 71° and 109°.\(^18\) It results that the structures can be constructed by placing the metal at a center of a cube and the four ligands at four vertices such that no ligands are trans (Figure 1). This gives three possible structures, which have the shape of a tetrahedron, \(\text{td}\), an umbrella, \(\text{umb}\), and a 4-legged-piano stool, \(\text{pst}\), as shown in the top part of Figure 2. In \(\text{td}\), all coordination sites are equivalent and all bond angles are 109°. In \(\text{umb}\), there are three basal hydrides, \(\text{H}_b\), and one apical hydride, \(\text{H}_a\), with \(\text{H}_b–\text{M}–\text{H}_a\) and \(\text{H}_a–\text{M}–\text{H}_b\) angles of 71° and 109°, respectively. In \(\text{pst}\), the four hydrides are equivalent, and the cis \(\text{H}–\text{M}–\text{H}\) angle is equal to 71°. Changing the metal and

![Figure 1](image-url)
the ligands lead to different angles, which influence the opening of the umbrella and the height of the piano stool.

When one of the three hydrogen atoms is replaced by an X ligand, tetrahedral, umbrella, and 4-legged piano stool isomeric structures are still found (Figure 2 bottom). Only the umbrella complex has nonequivalent sites. The presence of X gives rise to two isomers, namely, umb, where X is at the apical site and umbb, where X is at the basal site.

**Validation of the Use of Single Reference Methods for**

**d4 MH3X.** The geometries of MH3X were optimized at a DFT level and the energetics were obtained with methods 2, 3, and 4, DFT (PBE0 and RSH+MP2), and CCSD(T), respectively. However, the use of single reference (SR) methods for the description of unsaturated species could be questionable. Even with truncation of excitations at the double and triple level, single reference methods, such as CCSD(T), recover a large portion of the dynamic correlation energy, but may fail when a significant amount of nondynamical correlation energy is present. In contrast, multireference methods, such as MRCI or MR-ACPF, recover both dynamical and nondynamical correlation energies.

To check the validity of the SR calculations for d4 MH3X, single points MRCI and MR-ACPF-2a calculations were carried out for the three isomers of RuH4 (td, umb, and pst). For these MRCI and MR-ACPF-2a calculations, the reference space was built from valence CASCCF orbitals (12 electrons, 10 orbitals) and the effect of the rejected configurations was estimated using second order perturbation theory (Table 1).

All calculations show that the umb and pst isomers are lower in energy than the td isomer. Furthermore, all methods give a similar energetic pattern. Umb and pst are lower than td by 9.1 to 12.9 and 9.4 to 14.7 kcal mol\(^{-1}\), using the PBE0, RSH+MP2, and CCSD(T) methods, respectively. There is a marginal tendency for pst to have a lower energy than umb. It is also worth noting that the range separated hybrid RSH+MP2 calculations and the single reference coupled cluster calculations give values that are close. This overall good agreement between the single reference and multireference methods suggests that these unsaturated metal species do not contain a significant amount of nondynamical correlation energy. This indicates that a single reference method can be used for the computational study of d4 MH3X.

In addition, several diagnostics have been suggested to evaluate the SR/MR character for molecular systems. The T1 and D1 diagnostics\(^{33}\) (the Frobenius norm and matrix 2-norm of coupled cluster amplitudes for single excitations respectively) are arguably the most widely used diagnostics for SR coupled cluster calculations. Alternatively, the weight (Cfi\(^{2}\)) of the leading configuration state function in a MRCI can be used to determine the multireference character. The T1 and D1 diagnostics were tested primarily on small organic molecules, and it was suggested that any T1 diagnostic larger than 0.02 and any D1 diagnostic larger than 0.05 invalidate the use of SR methods. However, a recent study of Jiang et al. pointed out that the criteria established for organic molecules (T1 < 0.02 and D1 < 0.05) are no longer valid for molecular systems with transition metals.\(^{34}\) On the basis of a statistical analysis of a set of 225 species, these authors proposed T1 < 0.05 and D1 < 0.15 as amplitude criteria for the validation of SR-based methods in the case of d-block energetics. The T1 and D1 diagnostics of the single reference CCSD(T) calculations are reported in Table 2 for the three isomers of RuH4 (td, umb, and pst).

| Table 2. T1 and D1 Diagnostics of the Single Reference CCSD(T) Calculations and Weight of the Leading Configuration (Cfi\(^{2}\)) of MRCI and MRACPF2 Calculations for the Three Isomers of RuH4 |
|---|---|---|
| td | umb | pst |
| T1 | 0.048 | 0.045 | 0.043 |
| D1 | 0.091 | 0.100 | 0.089 |
| Cfi\(^{2}\) (MRCI) | 0.82 | 0.84 | 0.84 |
| Cfi\(^{2}\) (MRACPF2) | 0.81 | 0.83 | 0.83 |

and pst; the weight of the leading configuration (Cfi\(^{2}\)) from the multireference MRCI or MR-ACPF-2a calculation is also indicated. For the three isomers of RuH4, both T1 and D1 diagnostics from SR CCSD(T) calculations are less than the values suggested by Jiang et al. as acceptable limits. In addition, the weight of the leading configuration is higher than 0.80 for the MRCI and MR-ACPF-2a calculations. This confirms that these unsaturated molecular systems do not contain any substantial MR character. Single reference methods are thus validated for the calculations of RuH4. This validation applies to all d4 MH3X and MX4 studied in this article (T1 and D1 diagnostics given in the Supporting Information for RuH3X and OsH3X).

RuH3X. (a). RuH4. The DFT calculations show that RuH4 has td, umb, and pst minima. The umb and pst isomers have similar energies and both species are more stable than the tetrahedral structure by about 9 kcal mol\(^{-1}\). These results differ from those obtained by Landis. Molecular mechanic (MM) calculations including hybridization and resonance in a valence bond (VB) approach (Hypervalent-VALBOND (HV-VB) MM
method)\textsuperscript{9} give the same three minima but show that \textit{td} is the most stable isomer. DFT calculations using a different ECP and a different functional for the metal\textsuperscript{7a,9} from the one used in this work give \textit{td} and \textit{pst} as the only minima, with \textit{pst} being 8.6 kcal mol\textsuperscript{-1} lower in energy relative to \textit{td}. The Hartree–Fock and MP2 calculations by Hall et al. show several isomers with relative energies that depend on the computational method used.\textsuperscript{10}

The essential structural features of the three isomers of \textit{RuH}_4 are shown in Figure 3. The \textit{Ru}–\textit{H} bond is 1.55 Å in \textit{td}. In \textit{umb}, the apical \textit{Ru}–\textit{H} bond of 1.51 Å is slightly shorter than the basal \textit{Ru}–\textit{H} bond of 1.54 Å. The \textit{H}_a–\textit{Ru}–\textit{H}_b and the \textit{H}_b–\textit{Ru}–\textit{H}_a angles of 64° and 102°, respectively, are close to the corresponding ideal values of 71° and 109°. The acute \textit{H}_b–\textit{Ru}–\textit{H}_a angle leads to a nonbonded \textit{H}···\textit{H} distance of 1.62 Å, which is less than the sum of the van der Waals (vdW) radii of the two hydrogens (1.20 Å).\textsuperscript{35} In \textit{pst}, the \textit{Ru}–\textit{H} bond of 1.53 Å and the angles of 67° and 103° for cis \textit{H}–\textit{Ru}–\textit{H} and trans \textit{H}–\textit{Ru}–\textit{H} angles, respectively, are similar to the ideal values of 71° and 109°. In \textit{pst}, the shorter nonbonded \textit{H}···\textit{H} distance is 1.69 Å, which is less than the sum of the vDW radii of the two hydrogens although slightly longer than in \textit{umb}. Therefore, both \textit{umb} and \textit{pst} have short nonbonded \textit{H}···\textit{H} distances that are not present in the \textit{td} structure wherein the \textit{H}···\textit{H} nonbonded distance is 2.54 Å.

The transition states between these minima were searched using method 2 to evaluate the depth of the associated wells on the potential energy surface and the possibility of exchange between the three minima (Figure 3). A transition state of 8.7 kcal mol\textsuperscript{-1} above \textit{td} was located between \textit{td} and \textit{umb} and another transition state between \textit{umb} and \textit{pst} with an energy of 2.9 kcal mol\textsuperscript{-1} above \textit{umb}. No transition state was identified between \textit{td} and \textit{pst}. The tetrahedral structure is thus separated from the other minima by a significant barrier, but there is a low energy barrier between \textit{umb} and \textit{pst}. The structural features of the two transition states are shown in Figure 3. The transition state between \textit{td} and \textit{umb} isomers inverts the configuration at the metal as it inverts the umbrella. At the transition state, the ruthenium atom is in the equatorial plane formed by the three \textit{H}_a hydrogens and the \textit{Ru}–\textit{H}_b bond is perpendicular to the equatorial plane; this inversion at the metal occurs without significant change in the \textit{Ru}–\textit{H} bond distances. The transformation of \textit{umb} to \textit{pst} is a concerted swinging motion of two \textit{H}_a relative to the plane defined by \textit{Ru}–\textit{H}_b and the third \textit{Ru}–\textit{H}_b bonds. At the transition state, \textit{Ru}–\textit{H}_a and two \textit{Ru}–\textit{H} bonds are coplanar. This transformation also occurs without significant change of the \textit{Ru}–\textit{H} bond distances.

In \textit{RuH}_4F, the isomers \textit{td} and \textit{pst} are isoenergetic. The \textit{Ru}–\textit{H} bond lengths and the angles between the \textit{Ru}–\textit{H} bonds are similar to that of \textit{RuH}_4 in the corresponding isomers. The \textit{Ru}–\textit{F} is longer at the basal site of \textit{umb} than in \textit{td} (1.85 and 1.80 Å, respectively). The \textit{F}–\textit{Ru}–\textit{H} angle is nearly 92° in

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure3.png}
\caption{DFT optimized structures of minima and transition states of interconversion between minima for \textit{RuH}_4. The distances are in Å and the angles in degrees. The DFT energies are given in kcal mol\textsuperscript{-1} relative to the \textit{td} structure, using method 2.}
\end{figure}

\begin{figure}[h]
\centering
\includegraphics[width=\textwidth]{figure4.png}
\caption{DFT energies, in kcal mol\textsuperscript{-1}, for \textit{RuH}_4X (\textit{X} = F, CF\textsubscript{3}, CH\textsubscript{3}, H, SiF\textsubscript{3}, and SiH\textsubscript{3}) complexes, relative to the tetrahedral isomer, \textit{td}. The energies are in blue for the umbrella with apical \textit{X}, \textit{umb}, in red for the umbrella with basal \textit{X}, \textit{umb}\textsubscript{b}, and in green for the 4-legged piano-stool structure, \textit{pst}. Light green is used to indicate structures with hydride bridging \textit{Ru}–\textit{Si}.}
\end{figure}
systematically smaller than the corresponding C–Ru–H angles in the non-tetrahedral geometry. For the pst structure, one of the hydrogen of the SiH₃ group bridges the Ru–Si bond. The ability for the silyl group and notably the SiH₃ group to enter in the bridging situation is well documented and is not the focus of this study.⁶⁶

As a partial summary, these calculations show that all RuH₃X have several isomeric structures, among which the tetrahedral structure is not always the most stable isomer. For X = F and CH₃, the td structures are isoenergetic with the other isomers. In the case of X = H, CF₃, SiH₃ and SiF₃, there is a clear preference for non-tetrahedral structures especially with the silyl ligands. The structures of the td, umb, umb pst isomers are rather similar for all X and, thus, the structure of RuH₄ is a good representative of the entire set. There is a noticeable tendency for the X groups that favor significantly the non-tetrahedral structures (SiH₃ and SiF₃) to have a structure close to that of RuH₄. In contrast, for the X groups that do not favor the non-tetrahedral structures (F and CH₃), the X–Ru–H angle is systematically larger than the corresponding H–Ru–H angle. This structural pattern will reveal through-space interaction in the non-tetrahedral structures, as it will be discussed later.

For all systems, the various non-tetrahedral structures are separated by low energy barriers and therefore non-tetrahedral structures are fluxional. However, the non-tetrahedral structures are separated from the td structure by a barrier that is significantly higher than their difference in energy with the exception of RuH₃SiF₃ where the energy barrier is very close to the difference in energy. See Supporting Information for further details.

OsH₃X (X = H, F, CF₃, CH₃, SiH₃, and SiF₃). Three minima were located for OsH₄ (Figure 6). In contrast to the preference for non-tetrahedral structures found for RuH₄, the three isomers of OsH₄ are isoenergetic. Similar results have been obtained by Landis et al. with umb and pst being 3.5 and 5.2 kcal mol⁻¹ above the td.⁷⁷ The structures of td, umb, and pst are similar to those found for RuH₄. An energy barrier of 10.9 kcal mol⁻¹ relative to td separates td and umb,⁷⁷ and an energy barrier of 5.8 kcal mol⁻¹ relative to umb separates umb and pst.

The DFT energies for OsH₃X (X = F, CF₃, CH₃, H, SiF₃, and SiH₃), relative to the td isomers, are shown as a histogram in Figure 7 using the convention already used in Figure 4 for RuH₄X. There is a significant difference between the isoelectronic ruthenium and osmium complexes. For X = F, CH₃, and CF₃, the tetrahedral structure is more stable than any non-tetrahedral isomer. For X = H, the tetrahedral and non-tetrahedral isomers have similar energies. For X = SiH₃ and SiF₃, the preference for the non-tetrahedral structure remains, but the difference in energy between td and non-tetrahedral structures decreases by about 10 kcal mol⁻¹, compared to the ruthenium complexes. Thus, everything being equal, replacing ruthenium by osmium increases the stability of the tetrahedral isomer relative to the non-tetrahedral.

[ReH₄X]⁺, [RhH₄X]⁺, and [IrH₄X]⁺. To further understanding of the effect of the metal on the structural preference of d⁴ MH₄⁺ the isoelectronic anionic [ReH₄X]⁻ and cationic [RhH₄X]⁺, [IrH₄X]⁺ species were optimized with method 2, and single points RSH+MP2 energies were obtained with method 3.

[ReH₄]⁻ and [IrH₄]⁺ have three minima while [RhH₄]⁺ has two minima, Figure 8. The anionic [ReH₄]⁻ has a preference for the tetrahedral structure while the cationic [RhH₄]⁺ and

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Figure 5. RSH+MP2 (top) and CCSD(T) (bottom) energies, in kcal mol⁻¹, for RuH₄X (X = F, CF₃, CH₃, H, SiF₃, and SiH₃) complexes, relative to the tetrahedral isomer, td. The energies are in blue for the umbrella with apical X, umb, in red for the umbrella with basal X, umb, and in green for the 4-legged piano-stool structure, pst.

umb₃, which is notably different from the H₃–Ru–H₄ angle of 63°. For RuH₃CH₃ the td, umb₃ and pst have similar energies and umb is significantly higher in energy. Although the four minima have shapes that are overall similar to the corresponding minima of RuH₄, some structural aspects need to be mentioned. The Ru–C bond length of 1.95 Å in td is slightly shorter than in the other isomers (1.96 to 2.1 Å). In umb₃ the methyl group moves away from the axis of the umbrella as indicated by the CH₃–Ru–H₄ angle of 80°, which is larger than the H₃–Ru–H₄ angle of 66°. Likewise, in pst, the CH₃–Ru–H₄ angle of 81° is larger than that of 63° between two cis Ru–H bonds. RuH₃CF₃ also has four minima, the most stable ones being umb₃ and pst by about 8 kcal mol⁻¹ at the DFT level. The structural features of the RuH₃CF₃ and RuH₃CH₃ complexes are similar.

RuH₃SiH₃ and RuH₃SiF₃ also have the four minima like were found for RuH₃CH₃ and RuH₃CF₃. However, the energy pattern for the alkyl and the silyl ligands is significantly different. For the silyl ligand, all non-tetrahedral structures are at significantly lower energy than the tetrahedral structure by about 18 kcal mol⁻¹ at the DFT level. In addition, the influence on the geometries is also different for the silyl and the methyl ligands. For instance, in the umb₃ isomer of RuH₃SiH₃, the Ru–Si distance of 2.26 Å is slightly shorter than in td, 2.30 Å, and the Si–Ru–H angles are similar to the corresponding H–Ru–H angles. In umb₃ of RuH₃SiH₃ the Si–Ru–H₄ angle is equal to the H₃–Ru–H₄ angle. Thus, the Si–Ru–H angles are smaller than the H₃–Ru–H₄ angle in the non-tetrahedral geometry. For the pst structure, one of the hydrogen of the SiH₃ group bridges the Ru–Si bond. The ability for the silyl group and notably the SiH₃ group to enter in the bridging situation is well documented and is not the focus of this study.⁶⁶

As a partial summary, these calculations show that all RuH₃X have several isomeric structures, among which the tetrahedral structure is not always the most stable isomer. For X = F and CH₃, the td structures are isoeenergetic with the other isomers. In the case of X = H, CF₃, SiH₃ and SiF₃, there is a clear preference for non-tetrahedral structures especially with the silyl ligands. The structures of the td, umb, umb pst isomers are rather similar for all X and, thus, the structure of RuH₄ is a good representative of the entire set. There is a noticeable tendency for the X groups that favor significantly the non-tetrahedral structures (SiH₃ and SiF₃) to have a structure close to that of RuH₄. In contrast, for the X groups that do not favor the non-tetrahedral structures (F and CH₃), the X–Ru–H angle is systematically larger than the corresponding H–Ru–H angle. This structural pattern will reveal through-space interaction in the non-tetrahedral structures, as it will be discussed later.

For all systems, the various non-tetrahedral structures are separated by low energy barriers and therefore non-tetrahedral structures are fluxional. However, the non-tetrahedral structures are separated from the td structure by a barrier that is significantly higher than their difference in energy with the exception of RuH₃SiF₃ where the energy barrier is very close to the difference in energy. See Supporting Information for further details.

OsH₃X (X = H, F, CF₃, CH₃, SiH₃, and SiF₃). Three minima were located for OsH₄ (Figure 6). In contrast to the preference for non-tetrahedral structures found for RuH₄, the three isomers of OsH₄ are isoenergetic. Similar results have been obtained by Landis et al. with umb and pst being 3.5 and 5.2 kcal mol⁻¹ above the td.⁷⁷ The structures of td, umb, and pst are similar to those found for RuH₄. An energy barrier of 10.9 kcal mol⁻¹ relative to td separates td and umb,⁷⁷ and an energy barrier of 5.8 kcal mol⁻¹ relative to umb separates umb and pst.

The DFT energies for OsH₃X (X = F, CF₃, CH₃, H, SiF₃, and SiH₃), relative to the td isomers, are shown as a histogram in Figure 7 using the convention already used in Figure 4 for RuH₄X. There is a significant difference between the isoelectronic ruthenium and osmium complexes. For X = F, CH₃, and CF₃, the tetrahedral structure is more stable than any non-tetrahedral isomer. For X = H, the tetrahedral and non-tetrahedral isomers have similar energies. For X = SiH₃ and SiF₃, the preference for the non-tetrahedral structure remains, but the difference in energy between td and non-tetrahedral structures decreases by about 10 kcal mol⁻¹, compared to the ruthenium complexes. Thus, everything being equal, replacing ruthenium by osmium increases the stability of the tetrahedral isomer relative to the non-tetrahedral.

[ReH₄X]⁺, [RhH₄X]⁺, and [IrH₄X]⁺. To further understanding of the effect of the metal on the structural preference of d⁴ MH₄⁺ the isoelectronic anionic [ReH₄X]⁻ and cationic [RhH₄X]⁺, [IrH₄X]⁺ species were optimized with method 2, and single points RSH+MP2 energies were obtained with method 3.

[ReH₄]⁻ and [IrH₄]⁺ have three minima while [RhH₄]⁺ has two minima, Figure 8. The anionic [ReH₄]⁻ has a preference for the tetrahedral structure while the cationic [RhH₄]⁺ and
[IrH₄]⁺ have a preference for non-tetrahedral geometries. The preference for the non-tetrahedral relative to the tetrahedral structures is stronger for the cationic than for the neutral species; for [IrH₄]⁺ the preference for non-tetrahedral structure increases by over 25 kcal mol⁻¹ relative to RuH₄. For [IrH₄]⁺, the preference for the non-tetrahedral structure is more marked than for OsH₄ but less so than for RhH₄⁺. The influence of the total charge and the nature of metal on the preferential structures of d⁴[MH₄]⁺ (q = −1, 0, and 1) is thus important. Consequently, going from the left to the right of the periodic table increases the preference for non-tetrahedral geometries but going down a column decreases this preference.

The geometries of the minima for [ReH₄]⁻, [RhH₄]⁺, and [IrH₄]⁺ are similar to those previously found for RuH₄. The H–M–H angle shows a pattern that parallels the preference for the non-tetrahedral structures. In the umb isomer of [ReH₄]⁻, the H₁–M–H₂ and the H₃–M–H₄ angles are larger than in the neutral systems. A similar pattern is seen in the pst isomer. The opposite trend is found in the cationic systems. For [IrH₄]⁺, the minima are those of the type shown in Figure 2. In the case of [RhH₄]⁺, the pst isomer is a highly fluxional structure where the decrease of a cis H–Rh–H angle from 63 to 38° occurs with essentially no energy cost. No umb type was located as minimum. Overall, the structures of isoelectronic d⁴ MH₄ are thus sensitive to the total charge; a negative charge opens the umbrella and flattens the 4-legged piano stool and a positive charge closes the umbrella and raises the height of the 4-legged piano stool. Going down a column of the periodic table also opens the umbrella and flattens the 4-legged piano stool.

The histograms for the relative energies of the isomers for [ReH₃X]⁻, [RhH₃X]⁺, and [IrH₃X]⁺ are shown in the Supporting Information. The pattern of energies can be understood from the pattern obtained for the ruthenium and osmium complexes. For [ReH₃X]⁻ (X = F, CF₃, CH₃, SiF₃, and SiH₃) the preference for the tetrahedral structure decreases from CH₃ to SiF₃. However, the combined effect of the negative charge and the sd metal results in a general preference for the tetrahedral coordination for any X, in the case of [ReH₃F]⁻ the tetrahedral structure is the only minimum. For X = CH₃, CF₃, and SiH₃, the tetrahedral structures are more stable than the non-tetrahedral structures, and, in the case of [ReH₃SiF₃]⁻, pst and td are isoenergetic. For [IrH₃X]⁺ the preference for non-tetrahedral structures is found for X = SiF₃ and SiH₃. For X = CH₃ and CF₃, non-tetrahedral structures are preferred except for the umb structure with apical X group that is significantly less stable than td. In the case of [IrH₃F]⁺, the umb structure with equatorial F is marginally more stable than td, but the umb structure with apical F is significantly less stable than td. In the case of [RhH₃X]⁺, non-tetrahedral structures are more stable than td for X = H, SiF₃, and SiH₃. For X = CF₃ and CH₃, structures of different natures than found for...
all other metal complexes are found to be more stable than tetrahedral structure. No comparison can be done with the other cases, and this was not explored further.

RuX₄. The diamagnetic homoleptic Ru(C₆H₁₃)₄, and Ru(o-Tolyl)₄, Ru(Mesityl)₂, Os(SiH₃)₂, Os(o-Tolyl)₂, Ir(Mesityl)₂, which are the only 12-electron d⁴ MR₄ complexes to have been synthesized and fully characterized, are remarkably stable. They are colored as expected from the presence of low-lying empty metal d orbitals. In all cases, X-ray diffraction studies show that these complexes have a tetrahedral structure in the solid state. The NMR study of the tetramesityl complex of ruthenium shows a fluxional behavior attributed to the rotation of the ligand about the Ru–C bond. To better compare with the experimental structures, RuX₄ complexes were calculated for X = F, CH₃, CF₃, SiH₃, and SiF₃. The optimization of RuF₄ and Ru(CF₃)₄ yields only tetrahedral isomers. Two minima are found for Ru(CH₃)₄ with the td structure being 32.2 kcal mol⁻¹ below the pst structure. A preference for non-tetrahedral isomers is obtained with the silyl substituted complexes. For Ru(SiH₃)₄ and Ru(SiF₃)₄ the umb isomer is 12.6 kcal mol⁻¹ and 17.1 kcal mol⁻¹ more stable than the td isomer, respectively. In addition, structures with a pst shape in which the SiH₃ and SiF₃ ligands have hydrides or fluoride bridging the Ru–Si bond are also located as minima lower than the td isomer (−29.8 and −23.5 kcal mol⁻¹ for Ru(SiH₃)₄ and Ru(SiF₃)₄, respectively), but they will not be considered. These calculations on homoleptic complexes show that non-tetrahedral geometries are possible even with ligands more bulky than H. Furthermore, the energy trend obtained for RuHₓX (Figure 4) is also present for RuX₄. Fluorine is the ligand that most disfavors the non-tetrahedral geometry in the RuHₓX series; consequently RuF₄ has only a tetrahedral structure. The CH₃ ligand disfavors slightly the non-tetrahedral structures; consequently a non-tetrahedral structure is found as a high-lying secondary minimum. Finally, SiH₃ and SiF₃ ligands, which most favor non-tetrahedral structures in the RuHₓX series, lead to homoleptic complexes with preference for non-tetrahedral structures. The only exception to this trend is the case of Ru(CF₃)₄ whose tetrahedral structure cannot be predicted from the influence of a single CF₃ ligand on the structural preference of RuH₄(CF₃). Clearly, the structures of the alky and aryl complexes that have been synthesized do not display the diversity of possible structures for d⁴ RuX₄.

**DISCUSSION**

The bonding analysis of Landis and Weinhold as well as the ORSAM analysis of Hall show that d⁴ MH₄X and MX₄ have several structures. In this analysis, the metal uses the (n + 1)s and the nd orbitals, which are not occupied by the electrons of the metal lone pairs to establish the covalent metal–ligand bond. For an sd⁴ hybridization that applies to d⁴ tetracordinated complexes, the tetrahedral structure is one of the possible structures. The non-tetrahedral structures have umbrella and 4-legged piano stool shapes. However, the valence-bond analysis of Landis and Weinhold and the ORSAM model of Hall do not provide any information on the relative energies of the various isomeric forms. The DFT and ab initio calculations show that the metal and the ligands influence significantly the relative energies of the several isomers. In MH₄X, the non-tetrahedral structures have nonbonded distances between proximate atoms that are relatively short. This is the case in particular for MH₄ wherein the nonbonding H···H distances in the non-tetrahedral geometries are shorter than the sum of the vdW distances. The increased steric hindrance resulting from the replacement of the hydrogens in RuH₄ by relatively bulky alkyls and aryl ligands could have been the reason for all experimentally known systems to be tetrahedral. However, the calculations show that CH₃ and SiH₃ ligands lead to opposite structural preference, the former ligand increasing the preference for the tetrahedral structure and the latter for the non-tetrahedral structures. These trends apply to RuHₓX and RuX₄, showing that the steric effects of groups larger than hydrides do not determine the structural preferences. Calculations show that the metal also plays an important role on the structural preference. Going from the left to the right of the periodic table, that is, from anionic to cationic isostructural complexes increases the preference for the non-tetrahedral structures while going down a column of the periodic table increases the preference for the tetrahedral structures. Factors that also play a role are, among others, the relative electronegativity of the metal and the electron-donating/ electron-withdrawing ability of the ligands. To understand better these factors, a molecular orbital analysis, based on EHT calculations, has been carried out for RuH₄ and a NBO charge analysis complemented by a NRT analysis based on the calculated DFT densities have been carried out for RuH₄ and OsH₄.

(a) Molecular Orbital Analysis. The EHT total energies of the td, umb, and pst structures for RuH₄ give a preference for a non-tetrahedral structure. Increasing the electron-withdrawing ability of the ligand by lowering the hydrogen H₂ Coulombic integral gives a preference for a tetrahedral structure. This modeling of ligands reproduces the results of the DFT calculations. For instance (i) F and CH₃ are more electron-withdrawing than H and SiH₃ and thus favor the tetrahedral structure, (ii) Os 5d orbitals are higher in energy than Ru 4d orbitals and thus ligands appear to be more electron-withdrawing relative to Os than Ru. Likewise, in a cationic complex where the metal is more electron-attracting, all ligands appear more electron-donating. The reverse is true for anionic complexes.

A Walsh diagram is used to gain further insight into why EHT calculations reproduce qualitatively the structural preference obtained with higher-level calculations. Only the occupied molecular orbitals of RuH₄ are necessary in this analysis (Figure 9). The Cₓ or Cₓ group notations are used for labeling the molecular orbitals for the three structural forms. The td structure has two nonbonding d orbitals of e symmetry to host four electrons. The four Ru–H bonds are represented by two molecular orbitals of a₃ symmetry and two orbitals of e symmetry. Only the molecular orbitals of a₁ symmetry have a contribution on the hydrogen located on the z axis. In umb, all molecular orbitals of e symmetry have the same energy as in td because they have no contribution on the hydrogen on the z axis. The only molecular orbitals, which have different energies for td and umb, are the molecular orbitals of a₃ symmetry. They are constructed in a similar manner in td and umb. The lower one, 1a₁, is the in-phase combination of the metal 5s orbital and the 1s orbital of the four hydrogens. The higher orbital of a₁ symmetry, 2a₁, is mostly made of the in-phase combination of a d₃ with the four hydrogens. The contribution of the 5p orbitals has been found to be negligible in all molecular orbitals even if permitted by symmetry; it will not be mentioned further. The key point of the analysis is that, in 2a₁, the coefficient of the hydrogen on the z axis has the sign opposite to that of the three other hydrogens, which are in-phase with the torus part of d₃z.
The energy of 1a₁ is thus lower in umb than in td because the in-phase relationship between all hydrogens stabilizes the orbital when the hydrogen atoms are closer (umb vs td). In contrast, the energy of 2a₁ is higher in umb than in td because the out-of-phase relationship between the apical and basal hydrogens raises the energy of this molecular orbital when the distances between the nonbonded hydrogens are shorter and the overlap increases. Two antagonist effects are thus at work and umb is more stable than td if the energy lowering of 1a₁ is more important than the energy rise of 2a₁. The 1a₁ orbital is more located on the hydrogens than on the metal because the energy level H₄ of the hydrogen 1s orbital is always lower than the energy value H₄ of the 5s of the metal. In contrast, the 2a₁ orbital is more located on the metal when the energy of the hydrogen 1s orbital is higher than that of the metal d orbital, and it is more located on the hydrogen atoms when it is lower. Therefore, the energy rise of the 2a₁ orbital upon going from td to umb is small in the first case and large in the second case. Consequently, electron-donating ligands favor umb and electron-withdrawing ligands favor td.

A similar reasoning applies to the comparison between td and pst. In pst, the four nonbonding metal electrons are hosted in 2a₁ (dₓ²) and b₁ (dᵧ) orbitals (Figure 9). The four Ru–H bonding orbitals are described by an 1a₁ orbital made mostly from the metal s orbital, two orbitals of e symmetry made mostly from the metal dₓz and dᵧ orbitals, and a b₁ orbital made from dₓz−y². The contribution of the 5p orbitals has been found to be negligible even if permitted by symmetry, as already noted above. The 1a₁ orbital favors the pst structure because of the in-phase relationship between all hydrogens. The b₁ orbital disfavors the pst structure because of the out-of-phase relationship between the hydrogens. The two degenerate orbitals of e symmetry have essentially no influence on the structural preference because of the long distance between the two trans hydrogens. The influence of the energy level H₄ of the hydrogens on the variation of energy of the 1a₁ and b₁ orbitals between td and pst is similar to that obtained for the case of td vs umb. Electron-donating ligands favor pst and electron-withdrawing ligands favor td.

(b). NBO Charge Analysis. The NBO charges for RuH₄ and OsH₄ are shown in Figure 10. For these two species in which the metal is formally at a high oxidation state, M⁴⁺, the NBO charge at the metal is negative indicating, as currently recognized, that the calculated charge has no relation with formal oxidation state. The negative charge is small for the td complex but increases to a maximum value of −0.3 for the non-tetrahedral species. Accordingly, the charge on the hydrogens is very small and positive in RuH₄. In the non-tetrahedral structures, the hydrogens carry a positive charge; the largest positive charge is carried by the apical hydrogen of umb. The results are qualitatively similar for OsH₄, but the electronic density on the metal is smaller for Os than for Ru and the charge is even marginally positively charge in td. The charge distribution in the Ru and Os complexes does not give a rationale for the stronger energetic preference for non-tetrahedral structures in the case of Ru. However, the charge distribution indicates that electron-donor ligands are beneficial for the non-tetrahedral structures.

(c). NRT Analysis. A natural resonance (NRT) analysis of the DFT density was carried out for the three isomers, td, umb, and pst of RuH₄ and OsH₄. This analysis is aimed at determining the valence structures that describe the total electron density of these complexes in the various isomeric forms. This analysis has met some criticism but has been found to be qualitatively useful for understanding electronic structures. The results are shown in Figure 11. In the td isomer, the dominant structure has four covalent M−H bonds and two lone pairs at the metal. This valence structure accounts for 97% of the density for RuH₄ and 98% for OsH₄. For the umb isomer, the same structure accounts for only 80% of the total density for RuH₄ and 90% for OsH₄. The remaining part of the electron density is represented by structures with two covalent M−H bonds, three lone pairs at the metal, and a proton and a hydride characterizing an ionic interaction between the metal fragment and the hydrogen atoms. The structure with a proton at the apical site has higher weight (12% for RuH₄ and 6% for OsH₄) than that with a hydride at the apical site (7% for RuH₄ and 3% for OsH₄). The structure with the proton and the hydride at the basal sites has the smallest weight 1%. For the pst isomer, the main structure is that with four M−H bonds and two lone pairs at the metal (79% of the total electron density for RuH₄ and 89% for OsH₄). The remaining part of the electron density is represented by a valence structure with two covalent M−H bonds, three lone pairs at the metal, and a proton and a hydride on cis hydrogen.

Figure 9. Walsh diagram for the td to umb and pst transformation. The green orbitals are stabilized from td to either umb or pst and the red orbitals are destabilized. The green and red wriggling lines indicate positive and negative overlaps, respectively.

Figure 10. NBO charges in RuH₄ and OsH₄.
The density for OsH₄.

The power of the td (see Supporting Information). In all cases, the tetrahedral analysis was carried out for RuH₃X (X = F, CH₃, and SiH₃) contributes to the stability of the distances. This gives rise to an electrostatic interaction that opposite charges are on atoms that are at relatively short with the metal and also together since the two species with the ionic contributions, the proton and the hydride interact the contributions is 100%. In the valence structures displaying the umb structure that contributes to the total density. The values indicated are summed over all equivalent sites. It accounts for 21% of the density for RuH₄ and 11% of the density for OsH₄.

The total electron density is well described by these valence structures for all isomers of RuH₄ and OsH₄ since the sum of the contributions is 100%. In the valence structures displaying the ionic contributions, the proton and the hydride interact with the metal and also together since the two species with opposite charges are on atoms that are at relatively short distances. This gives rise to an electrostatic interaction that contributes to the stability of the umb and pst isomers. The structures with these ionic components have a higher weight for RuH₄ than for OsH₄, which could rationalize the preference for the umb and pst isomers in the case of RuH₄. The same analysis was carried out for RuH₃X (X = F, CH₃, and SiH₃) (see Supporting Information). In all cases, the tetrahedral structure is described by essentially a single structure with three M−H bonds, one M−X bond, and two lone pairs on the metal. The electron densities of umb and pst isomers are still dominated by the same valence structures but structures with two covalent bonds to the ruthenium, three lone pairs on the metal, and a positively charged group (H or X) and a negatively charged group (H or X) have relatively high percentages. It has not been possible to establish a quantitative correlation between the percentage of the valence structures with zwitterionic contributions and the relative energies of the non-tetrahedral and tetrahedral structures. However, it is likely that they play a key role in stabilizing the non-tetrahedral structures.

(d). Through-Space Interactions in the non-tetrahedral Structures from MO and NRT Analyses. The Walsh diagram and the NRT analyses provide a consistent interpretation of the stability of the non-tetrahedral structures of d⁴ RuH₄ although they use different languages and properties of the electronic structures. In the molecular orbital analysis, it was shown that the out-of-phase contribution between nearby “hydrogen” orbitals, which disfavors the umb and pst isomers relative to the td isomers, depends on the electron-donating power of the “hydrogen” ligands to the metal center. Ligands that are good electron donor contribute less to this orbital than ligands that are less good electron donor. When this destabilizing interaction is low, the through-space stabilizing interaction arising from the all in-phase low-lying occupied orbital dominates. Structures with the shortest distances between nonbonded atoms, that is, umb and pst, are favored relative to td. These stabilizing interactions do not create any covalent bonds between these hydrides but favor interatomic distances shorter than the sum of the vdW radii. The NRT analysis gives a complementary interpretation of the results. It confirms the absence of covalent bond between hydrides that are at distance shorter than the sum of the vdW radii (axial H₄ and basal H₄ in umb and two cis basal H in pst). However, the NRT establishes the presence of an electrostatic interaction between positively and negatively charged hydrides.

The analysis established for tetrahydride complexes can readily be generalized to complexes with any type of ligand. It provides an interpretation of the effect of the ligand using only the electron- withdrawing/electron-donating property of the ligands, but other factors can also contribute. For instance, fluorine is an electron-withdrawing ligand and thus the non-tetrahedral structures are disfavored. This result is suggested independently from the presence of the lone pairs on the fluorine lone pairs that would disfavor the presence of any group at short distance from fluorine as in the non-tetrahedral structures. Likewise, an alkyl group is reasonably electron- withdrawing and thus disfavors a non-tetrahedral structure. Steric effects between the alkyl group and the other ligands also disfavor the non-tetrahedral structures. At the other extreme, a good electron-donating group like a silyl group favors non-tetrahedral structures. In the case of a complex like RuH₄(SiH₄) the stabilizing interactions that have been described above are supplemented by the well-known through space interaction between a hydride and a nearby silyl group, originating from the ability of silicon to become hypervalent. This interaction, which has been known as SISHA⁴⁰ or IHI,⁴¹ is magnified when the silyl group is substituted by halide. This accounts for the preference for non-tetrahedral structures for RuH₄(SiH₄) and RuH₄(SiF₄). The fact that Ru(SiH₄)₂ and Ru(SiF₄)₂ also prefer non-tetrahedral structures show that interactions different from SISHA and IHI are at work.

CONCLUSIONS

DFT (PBE0 and range separated DFT, RSH+MP2) and CCSD(T) calculations have been used to determine the preferred structures of MH₃X³ (M = Ru, Os, Rh⁺, Ir⁺, and Re⁺; X = H, F, CH₃, CF₃, SiH₃, and SiF₃) and MX₂ (M = Ru; X = H, F, CH₃, CF₃, SiH₃, and SiF₃). Tetrahedral and non-tetrahedral structures are found as possible minima. The non-tetrahedral structures (umbrella and 4-legged piano stool shaped complexes) are those predicted for an sd³ hybridization following the analysis of valence bond of Landis et al. and the ORSAM analysis of Hall et al. The calculations show that non-tetrahedral structures can be energetically preferred to the expected tetrahedral geometry for certain metals and ligands. The non-tetrahedral isomers are more energetically preferred for 4d metal than for for 5d metals. Cationic complexes also tend to favor the non-tetrahedral isomers while anionic complexes do not. The relative energies of the non-tetrahedral and the tetrahedral structures are significantly influenced by the nature of the ligands. Electron-withdrawing ligands like halide and alkyl favor the conventional tetrahedral structure while ligands that are more electron donating like hydride and silyl.
favor the non-tetrahedral structures. The factors that determine these energy patterns have been discussed by means of a molecular orbital analysis based on EHT calculations and NBO charge analysis supplemented by an NRT analysis. All analyses show that electron-donating ligands favor the non-tetrahedral structures. Furthermore, the EHT and the NRT analyses suggest the presence of possible weak attractive interactions between atoms that are closer in umb and pst than in td. These weak attractive interactions are magnified when the ligands are good electron donor to the metal center.

## ASSOCIATED CONTENT

### Supporting Information
List of coordinates of all calculated structures with energies E and ZPE corrected energies (E + ZPE). Histograms of energies (methods 2 and 3) for [ReH$_3$X]$^+$, [RhH$_3$X]$^+$, and [IrH$_3$X]$^+$. Potential energy surfaces for the transformations between the various isomers of RuH$_6$X. NBO charges and NRT analysis for RuH$_6$X (X = F, CH$_3$, and SiH$_3$). This material is available free of charge via the Internet at http://pubs.acs.org.

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### Notes
The authors declare no competing financial interest.

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(39) The percentages shown in Figure 11 include the contribution of equivalent hydrogens.
